

## Crystal Structure

- Recall the Hamiltonian describing a system of ions, etc, that form a crystal. It has the full symmetry of 3D space (rigid motions: rotations, reflections, translations)
- Translation is particularly crucial:  $\vec{r} \rightarrow \vec{r} + \vec{a}$  is a symmetry for all vectors  $\vec{a}$ .
- A crystal is invariant only under discrete subset of translations:  $\{\mathcal{T}_{\vec{R}}\}$ , where  $\vec{R} = n_1\vec{a}_1 + n_2\vec{a}_2 + n_3\vec{a}_3$ .
- Example of spontaneous symmetry breaking: physical system has less symmetry than the Hamiltonian that describes it.
  - Happens for crystalline solids, but not for liquids & gases, which are homogeneous & isotropic.

# Mathematical Group Theory - Language of Symmetry

- Physicist's definition: A group is a set of symmetry operations, that is closed under composition of operations in the set.
  - Precise definition:
- A group is a set of elements  $\{g_i\}$ , and a rule for multiplying two  $g_i$ 's, with the following properties:
- (1)  $g_i g_j = g_k$ ;  $g_k$  is also a group element (closure)
  - (2) There is a special element  $e$  satisfying  $e g_i = g_i e = g_i$ . (Identity transformation)
  - (3) For every  $g_i$ , there is a unique  $g_i^{-1}$  satisfying  $g_i^{-1} g_i = g_i g_i^{-1} = e$ . (Inverse transformations.)
  - (4)  $(g_i g_j) g_k = g_i (g_j g_k)$ . (Associativity)

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NB: In general,  $g_i g_j \neq g_j g_i$  (Not commutative.)

Examples of groups: Rigid motions of space, continuous translations, discrete translations, rotations, ...

• Subgroup: A subgroup is a subset of some group, which is itself also a group. (In particular, it is closed under multiplication.)

• Example: Discrete translations  $\{T_{\vec{R}}\}$  ( $\vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3$ ) is a subgroup of the group of all translations.

• Symmetries / symmetry breaking: A larger symmetry group (e.g. continuous translations)

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- Precise statement of spontaneous symmetry breaking:
  - Hamiltonian has a symmetry described by some group  $G$ .
  - The system itself has a lower symmetry described by a group  $H$ , where  $H$  is a subgroup of  $G$ .

### Description of Crystal Structure

#### Bravais lattice:

- Set of points  $\{\vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3\}$ 
  - $n_i$ : integers
  - $\vec{a}_i$  must be linearly independent
  - $d=1$ :  $\{\vec{R} = n_1 \vec{a}_1\}$
  - $d=2$ :  $\{\vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2\}$
- Useful when  $\{\vec{R}\}$  are precisely the translation vectors that are crystal symmetries.

- Simplest monatomic ~~copper~~ crystals:  $\{\vec{R}\}$  ~~will~~ give the atomic positions.

### Example Bravais Lattices:

- Start with  $d=2$  (for ease of drawing)

(a) Square lattice "lattice constant"

$$\begin{array}{ccccccc} \cdot & \cdot & \cdot & & & & \\ \vec{a}_1 & \uparrow & \vec{a}_2 & = & \vec{a} \hat{x} & & \\ & \longrightarrow & & & & & \\ & \vec{a}_1 & & & & & \end{array}$$

(b) Hexagonal lattice (a.k.a. triangular lattice)

$$\begin{array}{ccccccc} \cdot & \cdot & \cdot & \cdot & & & \\ \cdot & \cdot & \cdot & \cdot & & & \\ \vec{a}_1 & \uparrow & \vec{a}_2 & = & \frac{a}{2} \hat{x} + \frac{\sqrt{3}}{2} a \hat{y} & & \\ & \longrightarrow & & & & & \\ & \vec{a}_1 & & & & & \end{array}$$

(c) Oblique lattice (no special symmetries — generic choice of  $\vec{a}_1$  &  $\vec{a}_2$ )

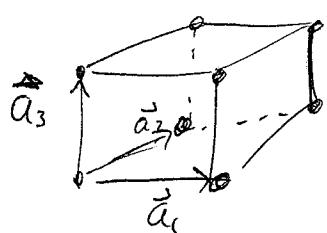
$$\begin{array}{ccccccc} \vec{a}_2 & \uparrow & \cdot & \cdot & \cdot & & \\ & \longrightarrow & & & & & \\ & \vec{a}_1 & & & & & \end{array}$$

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- $\vec{a}_1, \vec{a}_2 \rightarrow$  "primitive vectors"
- Choice of primitive vectors is not unique for a given lattice.  
(e.g.  $\vec{a}_1 = a\hat{x}$ ;  $\vec{a}_2 = a(\hat{x} + \hat{y})$  also works for square lattice)

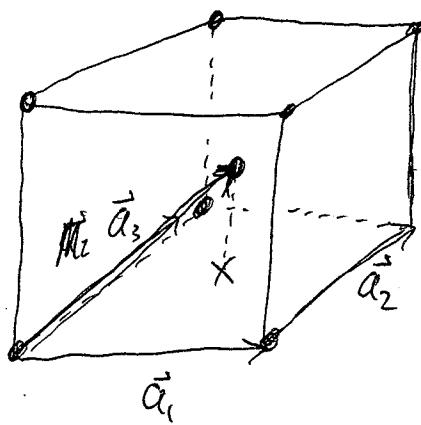
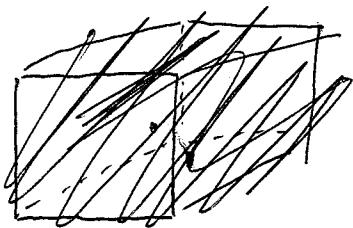
### d = 3 Lattices

#### (a) Simple cubic (example material: $P_0$ )



$$\begin{aligned}\vec{a}_1 &= a\hat{x} \\ \vec{a}_2 &= a\hat{y} \\ \vec{a}_3 &= a\hat{z}\end{aligned}$$

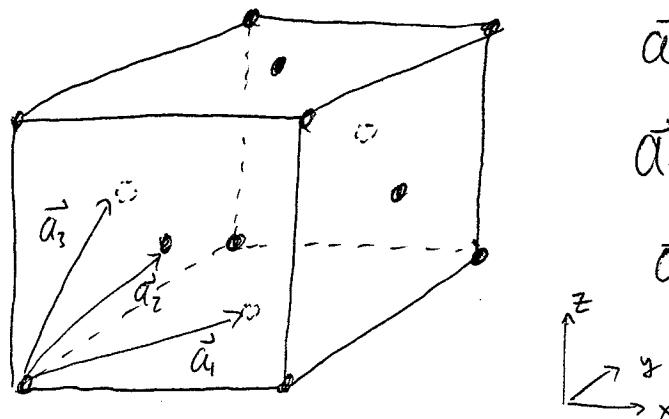
#### (b) BCC (ex: alkali metals $Li, Na, K, \dots$ ; Fe)



$$\begin{aligned}\vec{a}_1 &= a\hat{x} \\ \vec{a}_2 &= a\hat{y} \\ \vec{a}_3 &= \frac{a}{2}(\hat{x} + \hat{y} + \hat{z})\end{aligned}$$

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(c) FCC (ex: Cu, Ag, Au ; noble gases Ne, Ar, Kr, ... )



$$\vec{a}_1 = \frac{a}{2}(\hat{x} + \hat{y})$$

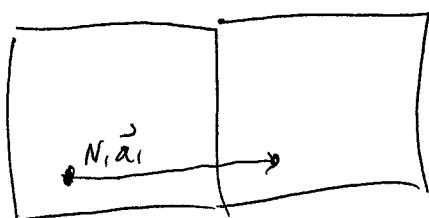
$$\vec{a}_2 = \frac{a}{2}(\hat{x} + \hat{z})$$

$$\vec{a}_3 = \frac{a}{2}(\hat{y} + \hat{z})$$

- Non-infinite crystals: Convenient to consider putting the crystal in a periodic box ("periodic boundary conditions"):

$$\left. \begin{array}{l} \{\vec{R}\}; \vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3 \\ \text{with } 0 \leq n_i \leq N_i - 1 \\ 0 \leq n_2 \leq N_2 - 1 \\ 0 \leq n_3 \leq N_3 - 1 \end{array} \right\} \begin{array}{l} N = N_1 N_2 N_3 \text{ total} \\ \text{lattice sites} \end{array}$$

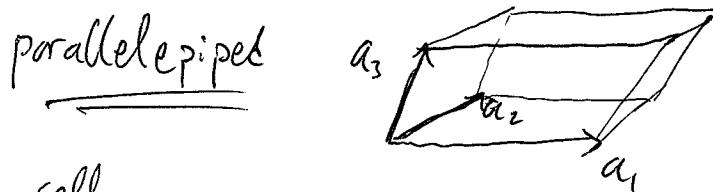
- Translation by  $N_1 \vec{a}_1$ ,  $N_2 \vec{a}_2$  or  $N_3 \vec{a}_3$  returns to same point.



## Unit Cells

- Primitive unit cell: Volume that fills space (with no overlaps) when translated by Bravais lattice vectors  $\{\vec{R}\}$ .

- example:  $\{x_1\vec{a}_1 + x_2\vec{a}_2 + x_3\vec{a}_3 \mid 0 \leq x_i \leq 1\}$



$$\text{cell volume} = |\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)|$$

- Infinite number of different possible primitive cells.  
(examples)
- All primitive unit cells have same volume.  
(since each contains exactly 1 lattice point.)
- Non-Primitive Unit Cell: Volume that fills space (no overlaps) when translated by some subset of  $\{\vec{R}\}$ .
  - example: Conventional cubic unit cells for BCC & FCC lattices

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- Wigner - Seitz Cell : Special primitive unit cell constructed by taking a given lattice point. The W-S cell is the volume enclosing that point where every point in the W-S cell is closer to the given lattice point than any other point.

ex: square:



The W-S cell has  
the full symmetry of  
the lattice.

hexagonal:



## More General Crystal Structures:

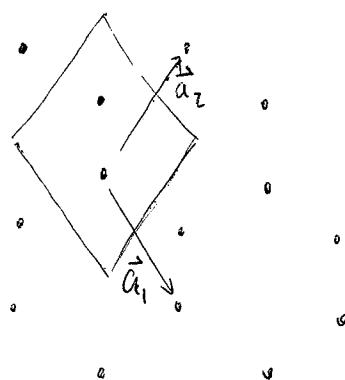
- Given a Bravais lattice, can put anything in a primitive unit cell:

?	?	?
?	?	?

### Lattice with Basis

- Basis:  $\{$  whatever we put in unit cell  
 $\}$  = some set of atoms.

### Example: Honeycomb lattice (graphene)



$\sim \vec{a}_1$  &  $\vec{a}_2$  are primitive vectors for hexagonal Bravais lattice

NB: All points of honeycomb lattice are equivalent by symmetry ... but not translation symmetry.

## Specifying

- To specify basis (of  $k$  atoms):

-  ~~$\vec{R}_i$~~ ;  $i=1, \dots, k$  give positions within unit cell  
 $\{\vec{c}_i\}$  Bravais lattice vectors  $\{\vec{R}\}$ .

After

- General atomic positions are  $\{\vec{R} + \vec{c}_i\}$

$\vec{c}_i$

- More examples (see crystal maker files online for images)

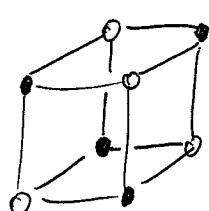
- Diamond lattice ( $C$ (diamond), Si, Ge)

- FCC with 2-site basis.

- Zincblende structure ( $GeAs$ )

- Like diamond, but 2 atoms in the basis are different.

- Sodium Chloride structure



simple cubic, with  $1/2$  atoms of one type,  $1/2$  of other type

(Can view as FCC with 2-site basis.)

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- Hexagonal Close Packed Lattice (HCP) : ( $Mg, Ti, Zn, \dots$ )

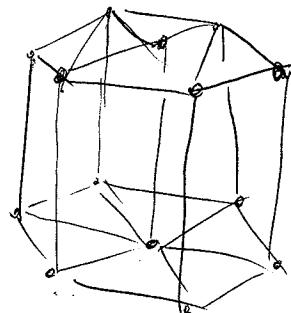
(3d)

- Based on hexagonal Bravais lattice, generated by :

$$\vec{a}_1 = a \hat{x}$$

$$\vec{a}_2 = a \left( \frac{1}{2} \hat{x} + \frac{\sqrt{3}}{2} \hat{y} \right)$$

$$\vec{a}_3 = c \hat{z}$$



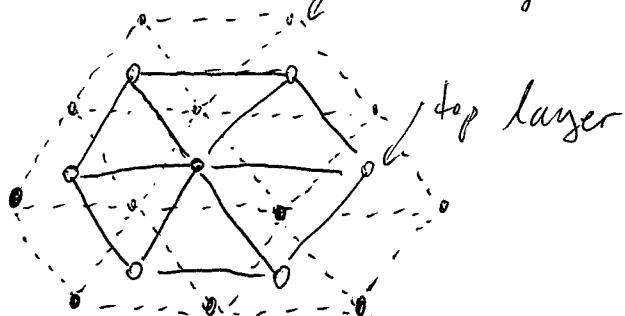
- Stacked triangular lattices.

- HCP is hexagonal Bravais lattice with basis

$$C_1 = 0$$

$$C_2 = \frac{\vec{a}_1}{3} + \frac{\vec{a}_2}{3} + \frac{\vec{a}_3}{2}$$

bottom layer



- Perovskite structure ( $BaTiO_3$ )

- based on simple cubic Bravais lattice